

Higher-order cochain interpolation

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June 29, 2010

Outline

Discrete representation of physical quantities

- Discretization: pointwise approach
- Discretization: geometric approach
- The de Rham map

Reconstruction of physical quantities

- The Whitney map
- Edge basis functions
- Finite volumes and Galerkin

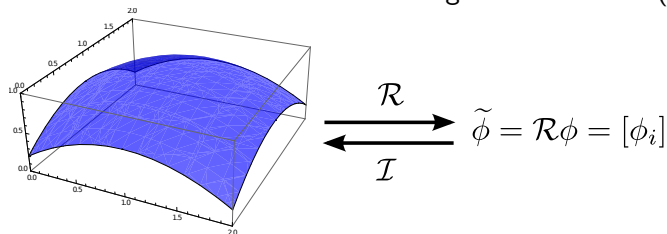
Conclusions

Discretization: pointwise approach

Problem? Reduce the number of degrees of freedom (dof)

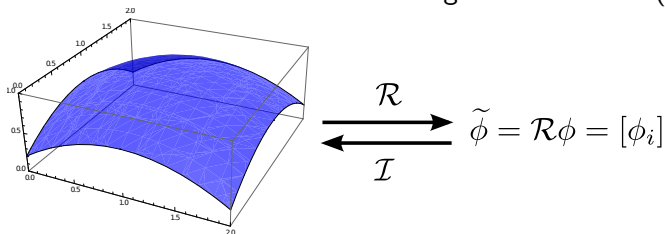
Discretization: pointwise approach

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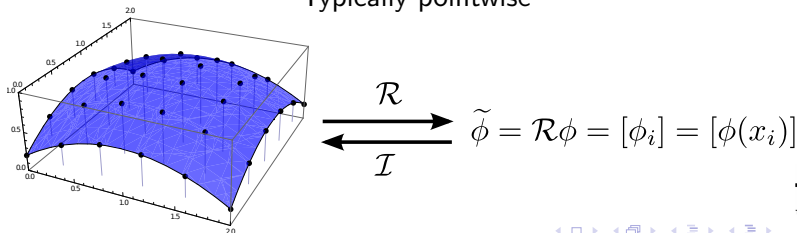


Discretization: pointwise approach

Problem? Reduce the number of degrees of freedom (dof)



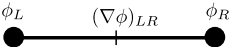
Typically pointwise

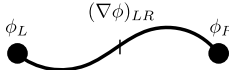


Differential operators become approximations

Differential operators become approximations

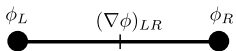
Pointwise discrete gradient (G):

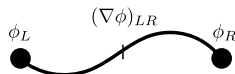

$$(\nabla\phi)_{LR} = \frac{1}{\Delta x}(\phi_R - \phi_L)$$


$$(\nabla\phi)_{LR} \approx \text{????}$$

Differential operators become approximations

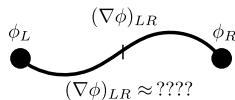
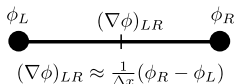
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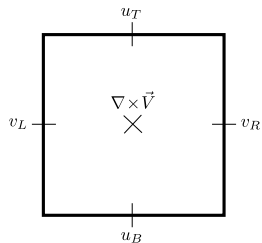

$$(\nabla\phi)_{LR} \approx \text{????}$$

Differential operators become approximations

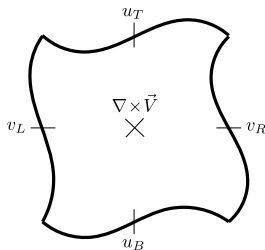
Pointwise discrete gradient (G):



Pointwise discrete curl (C):



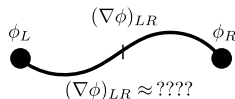
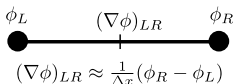
$$\nabla \times \vec{V} \approx \frac{1}{\Delta x}(v_R - v_L) + \frac{1}{\Delta y}(u_B - u_T)$$



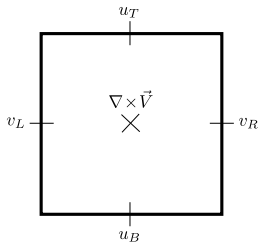
$$\nabla \times \vec{V} \approx ???$$

Differential operators become approximations

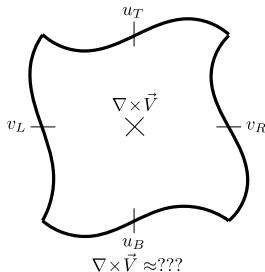
Pointwise discrete gradient (G):



Pointwise discrete curl (C):



$$\nabla \times \vec{V} \approx \frac{1}{\Delta x \Delta y} [(v_R - v_L)\Delta y + (u_B - u_T)\Delta x]$$



Discretization: geometric approach

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Fundamental theorem of calculus

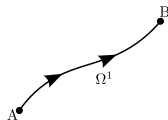
Kelvin-Stokes theorem

Gauss theorem

Discretization: geometric approach

Fundamental theorem of calculus

$$\int_{\Omega^1} \nabla \varphi \cdot d\vec{s} = \varphi(\Omega_B^1) - \varphi(\Omega_A^1)$$



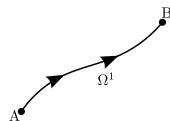
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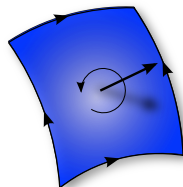
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Kelvin-Stokes theorem

$$\int_{\Omega^2} (\nabla \times \vec{V}) \cdot \vec{n} dS = \int_{\partial\Omega^2} \vec{V} \cdot d\vec{s}$$

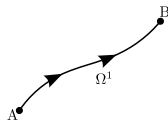


Gauss theorem

Discretization: geometric approach

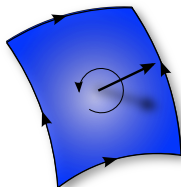
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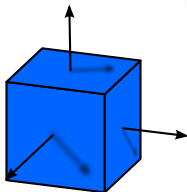
Kelvin-Stokes theorem

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Gauss theorem

$$\int_{\Omega^3} (\nabla \cdot \vec{V}) dV = \int_{\partial\Omega^3} \vec{V} \cdot \vec{n} dS$$



Differential geometry: Generalized Stokes Theorem

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$$\int_{\Omega^{k+1}} d\omega^k = \int_{\partial\Omega^{k+1}} \omega^k$$

With $\omega^k \in \Lambda^k$, space of continuous differential k -forms.

Differential geometry: Generalized Stokes Theorem

$$\int_{\Omega^{k+1}} d\omega^k = \int_{\partial\Omega^{k+1}} \omega^k$$

$$\langle d\omega^k, \Omega^{k+1} \rangle = \langle \omega^k, \partial\Omega^{k+1} \rangle$$

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$$k = 0 : \int_L d\phi^0 = \phi^0(B) - \phi^0(A) \quad d : \Lambda^0 \mapsto \Lambda^1 \quad (\text{maps points to lines})$$

$$k = 1 : \int_S d\alpha^1 = \int_{\partial S} \alpha^1 \quad d : \Lambda^1 \mapsto \Lambda^2 \quad (\text{maps lines to surfaces})$$

$$k = 2 : \int_V d\beta^2 = \int_{\partial V} \beta^2 \quad d : \Lambda^2 \mapsto \Lambda^3 \quad (\text{maps surfaces to volumes})$$

Differential geometry: Generalized Stokes Theorem

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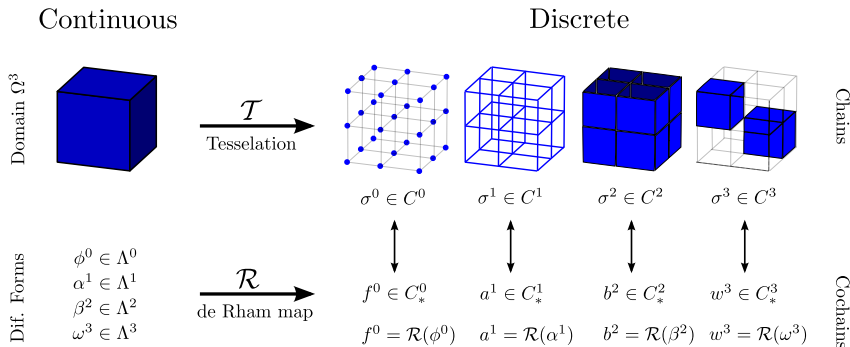
$$k = 2 : \int_V d\beta^2 = \int_{\partial V} \beta^2 \quad d : \Lambda^2 \mapsto \Lambda^3 \quad (\text{maps surfaces to volumes})$$

Constitutes an exact sequence (de Rham complex), since $dd = 0$

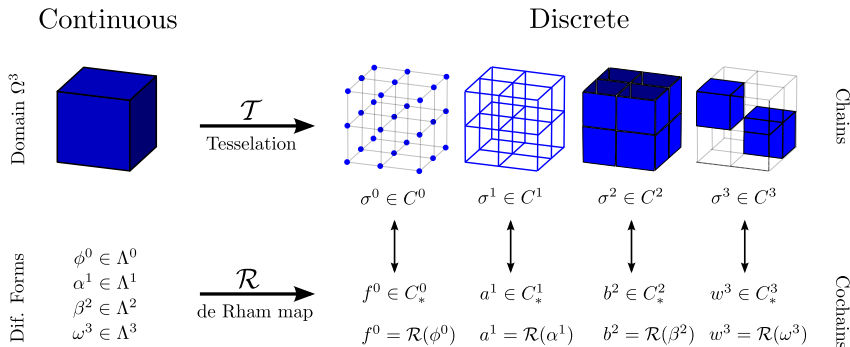
$$\mathbb{R} \hookrightarrow \Lambda^0 \xrightarrow[\nabla]{d} \Lambda^1 \xrightarrow[\nabla \times]{d} \Lambda^2 \xrightarrow[\nabla \cdot]{d} \Lambda^3 \xrightarrow{d} 0$$

Chains and cochains and the de Rham map

Chains and cochains and the de Rham map



Chains and cochains and the de Rham map



$$\mathcal{R} : \Lambda^k \mapsto C^k_*$$

$$\mathcal{R}(\alpha^k) = \left[\int_{\sigma_i^k \in \sigma^k} \alpha^k \right] = [a_i^k] = a^k$$

Exact differential operators

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$$\omega^{k+1} = d\alpha^k$$

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$$\omega^{k+1} = d\alpha^k$$
$$\mathcal{R}(\omega^{k+1}) = \mathcal{R}(d\alpha^k)$$

Exact differential operators

$$\begin{aligned}\omega^{k+1} &= d\alpha^k \\ \mathcal{R}(\omega^{k+1}) &= \mathcal{R}(d\alpha^k) \\ &= \left[\int_{\sigma_i^{k+1}} d\alpha^k \right]\end{aligned}$$

Exact differential operators

$$\begin{aligned}\omega^{k+1} &= d\alpha^k \\ \mathcal{R}(\omega^{k+1}) &= \mathcal{R}(d\alpha^k) \\ &= \left[\int_{\sigma_i^{k+1}} d\alpha^k \right] \\ &= \left[\int_{\partial\sigma_i^{k+1}} \alpha^k \right]\end{aligned}$$

Exact differential operators

$$\begin{aligned}\omega^{k+1} &= d\alpha^k \\ \mathcal{R}(\omega^{k+1}) &= \mathcal{R}(d\alpha^k) \\ &= \left[\int_{\sigma_i^{k+1}} d\alpha^k \right] \\ &= \left[\int_{\partial\sigma_i^{k+1}} \alpha^k \right] \\ &= \delta \left[\int_{\sigma_i^k} \alpha^k \right]\end{aligned}$$

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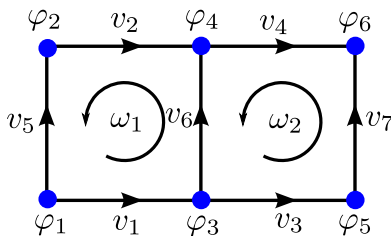
Exact differential operators

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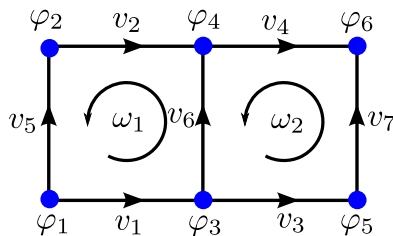
Exact differential operators

$$\begin{aligned}
 \omega^{k+1} &= d\alpha^k \\
 \mathcal{R}(\omega^{k+1}) &= \mathcal{R}(d\alpha^k) \\
 &= \left[\int_{\sigma_i^{k+1}} d\alpha^k \right] \\
 &= \left[\int_{\partial\sigma_i^{k+1}} \alpha^k \right] \\
 &= \delta \left[\int_{\sigma_i^k} \alpha^k \right] \\
 \mathcal{R}(\omega^{k+1}) &= \delta\mathcal{R}(\alpha^k) \\
 \Rightarrow \mathcal{R}d &= \delta\mathcal{R} \\
 \Rightarrow \delta\delta &= 0
 \end{aligned}$$

Exact differential operators



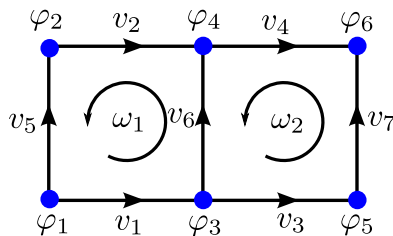
Exact differential operators



Gradient $\nu^1 = d\phi^0 \rightarrow v^1 = \delta f^0$

$$\begin{bmatrix} v_1 \\ v_2 \\ v_3 \\ v_4 \\ v_5 \\ v_6 \\ v_7 \end{bmatrix} = \begin{bmatrix} -1 & 0 & 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 & 0 & 1 \\ -1 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 & -1 & 1 \end{bmatrix} \begin{bmatrix} f_1 \\ f_2 \\ f_3 \\ f_4 \\ f_5 \\ f_6 \\ f_7 \\ f_8 \end{bmatrix}$$

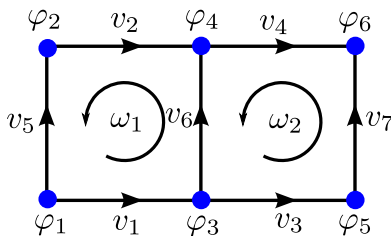
Exact differential operators



$$\text{Curl } \omega^2 = d\nu^1 \rightarrow w^2 = \delta v^1$$

$$\begin{bmatrix} w_1 \\ w_2 \end{bmatrix} = \begin{bmatrix} 1 & -1 & 0 & 0 & -1 & 1 & 0 \\ 0 & 0 & 1 & -1 & 0 & -1 & 1 \end{bmatrix} \begin{bmatrix} v_1 \\ v_2 \\ v_3 \\ v_4 \\ v_5 \\ v_6 \\ v_7 \end{bmatrix}$$

Exact differential operators



$$dd = 0 \rightarrow \delta\delta = 0$$

$$\begin{bmatrix} 1 & -1 & 0 & 0 & -1 & 1 & 0 \\ 0 & 0 & 1 & -1 & 0 & -1 & 1 \end{bmatrix} \begin{bmatrix} -1 & 0 & 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 & 0 & 1 \\ -1 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 & -1 & 1 \end{bmatrix} = 0$$

Exact differential operators

Discrete exact complex

$$\begin{array}{ccccccc}
 \mathbb{R} & \longrightarrow & \Lambda^0 & \xrightarrow{d} & \Lambda^1 & \xrightarrow{d} & \Lambda^2 & \xrightarrow{d} & \Lambda^3 & \longrightarrow & 0 \\
 & & \mathcal{R} \downarrow & & \mathcal{R} \downarrow & & \mathcal{R} \downarrow & & \mathcal{R} \downarrow & & \\
 \mathbb{R} & \longrightarrow & C_*^0 & \xrightarrow{\delta} & C_*^1 & \xrightarrow{\delta} & C_*^2 & \xrightarrow{\delta} & C_*^3 & \longrightarrow & 0
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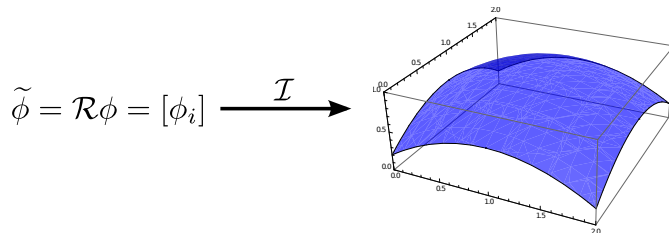
Cochain reconstruction: the Whitney map, \mathcal{I}

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Problem? Given a cochain, get a piecewise smooth reconstruction

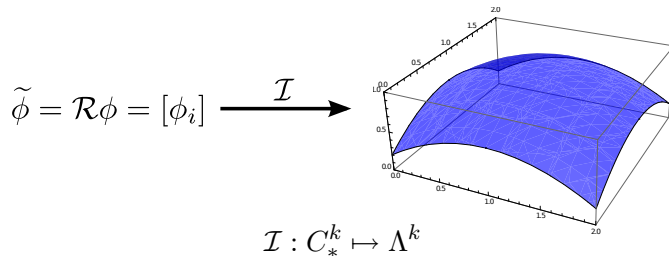
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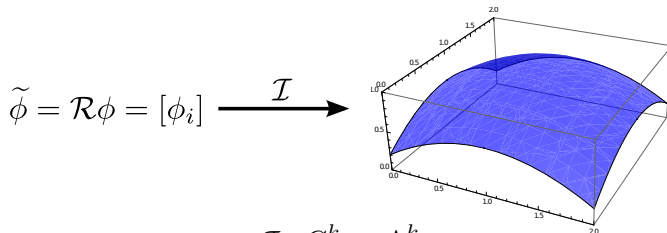
Problem? Given a cochain, get a piecewise smooth reconstruction



$$\mathcal{I}(a^k) = \sum_i a_i \epsilon_i^k(\mathbf{x}), \quad \epsilon_i^k(\mathbf{x}) \in \Lambda_p^k \subset \Lambda^k$$

Cochain reconstruction: the Whitney map, \mathcal{I}

Problem? Given a cochain, get a piecewise smooth reconstruction



$$\tilde{\phi} = \mathcal{R}\phi = [\phi_i] \xrightarrow{\mathcal{I}}$$

$$\mathcal{I} : C_*^k \mapsto \Lambda^k$$

$$\mathcal{I}(a^k) = \sum_i a_i \epsilon_i^k(\mathbf{x}), \quad \epsilon_i^k(\mathbf{x}) \in \Lambda_p^k \subset \Lambda^k$$

Such that:

$$\mathcal{R} \circ \mathcal{I} = \mathbb{1}, \quad \mathcal{I} \circ \mathcal{R} = \mathbb{1} + \mathcal{O}(h^p) \quad \text{and} \quad d\mathcal{I} = \mathcal{I}\delta$$

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$$\mathcal{R} \circ \mathcal{I} = \mathbb{1}$$

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$$\mathcal{R} \circ \mathcal{I}(a^k) = \mathcal{R}[\mathcal{I}(a^k)]$$

Cochain reconstruction: the Whitney map, \mathcal{I}

$$\mathcal{R} \circ \mathcal{I} = \mathbb{1}$$

$$\begin{aligned}\mathcal{R} \circ \mathcal{I}(a^k) &= \mathcal{R}[\mathcal{I}(a^k)] \\ &= \mathcal{R}\left[\sum_i a_i \epsilon_i^k(\mathbf{x})\right]\end{aligned}$$

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$$\Rightarrow \mathcal{R}[\epsilon_i^k] = \left[\int_{\sigma_j^k} \epsilon_i^k \right] = \delta_{ij} \quad (\text{Kronecker delta})$$

Cochain reconstruction: the Whitney map, \mathcal{I}

$$d\mathcal{I} = \mathcal{I}\delta$$

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$$d\mathcal{I}(a^k) = \mathcal{I}\delta(a^k)$$

Cochain reconstruction: the Whitney map, \mathcal{I}

$$d\mathcal{I} = \mathcal{I}\delta$$

$$\begin{aligned}d\mathcal{I}(a^k) &= \mathcal{I}\delta(a^k) \\d \sum_j a_j \epsilon_j^k &= \mathcal{I} \sum_j [\delta]_{ij} a_j\end{aligned}$$

Cochain reconstruction: the Whitney map, \mathcal{I}

$$d\mathcal{I} = \mathcal{I}\delta$$

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$$d \sum_j a_j \epsilon_j^k = \mathcal{I} \sum_j [\delta]_{ij} a_j$$

$$\sum_j a_j d\epsilon_j^k = \sum_i \sum_j [\delta]_{ij} a_j \epsilon_i^{k+1}$$

Cochain reconstruction: the Whitney map, \mathcal{I}

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$$\sum_j a_j d\epsilon_j^k = \sum_j a_j \sum_i [\delta]_{ij} \epsilon_i^{k+1}$$

$$\Rightarrow d\epsilon_j^k = \sum_i [\delta]_{ij} \epsilon_i^{k+1}$$

Exact de Rham complex

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Polynomial exact complex

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 \end{array}$$

Edge basis functions

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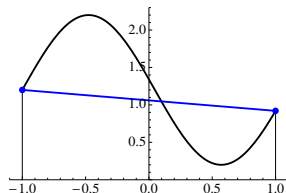
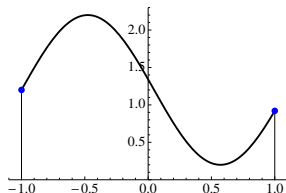
Nodal quantities: Lagrange interpolants

$$h_i^p(x_j) = \delta_{ij}$$

$$i, j = 1, \dots, p + 1$$

Edge basis functions

Nodal quantities: $p=1$



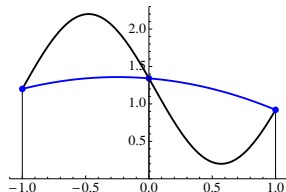
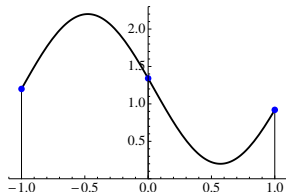
Reduce
(Discretize)
 \mathcal{R}

$$[\tilde{\varphi}_i] = [\varphi(-1), \varphi(1)]$$

Reconstruct
(Interpolate)
 \mathcal{I}

Edge basis functions

Nodal quantities: $p=2$



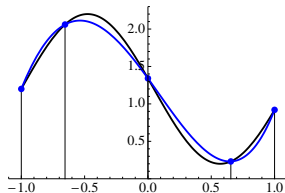
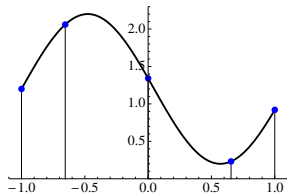
Reduce
(Discretize)
 \mathcal{R}

$$[\tilde{\varphi}_i] = [\varphi(-1), \varphi(0), \varphi(1)]$$

Reconstruct
(Interpolate)
 \mathcal{I}

Edge basis functions

Nodal quantities: $p=4$



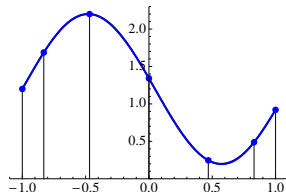
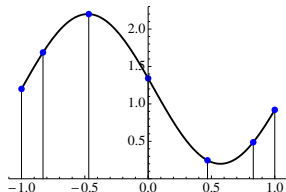
Reduce
(Discretize)
 \mathcal{R}

$$[\tilde{\varphi}_i] = [\varphi(-1), (\dots), \varphi(1)]$$

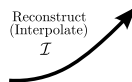
Reconstruct
(Interpolate)
 \mathcal{I}

Edge basis functions

Nodal quantities: $p=6$



$$[\tilde{\varphi}_i] = [\varphi(-1), (\dots), \varphi(1)]$$



Edge basis functions

Volume quantities: edge functions

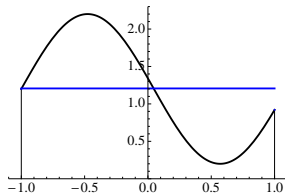
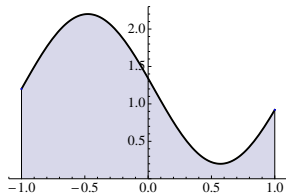
$$\int_{x_j}^{x_{j+1}} e_i^p(x) = \delta_{ij}$$

$$e_i^p(x) = - \sum_{n=1}^i \frac{d}{dx} h_n^{p+1}(x)$$

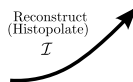
$$i, j = 1, \dots, p$$

Edge basis functions

Volume quantities: $p=0$

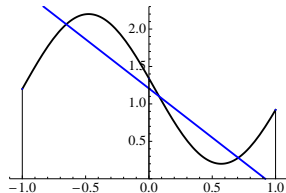
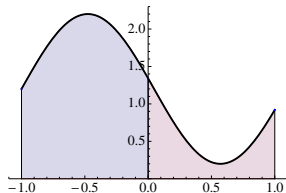


$$[\tilde{\varphi}_i] = [\bar{\varphi}_1, \bar{\varphi}_2]$$



Edge basis functions

Volume quantities: $p=1$



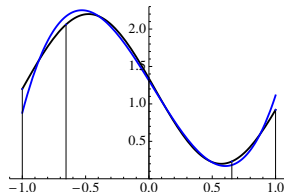
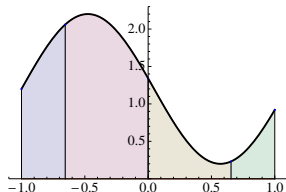
Reduce
(Discretize)
 \mathcal{R}

$$[\tilde{\varphi}_i] = [\bar{\varphi}_1, \bar{\varphi}_2, \bar{\varphi}_3]$$

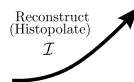
Reconstruct
(Histopolate)
 \mathcal{I}

Edge basis functions

Volume quantities: $p=3$

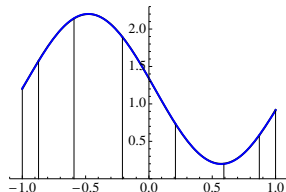
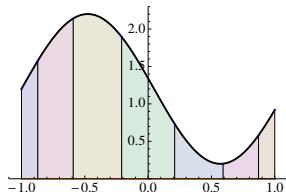


$$[\tilde{\varphi}_i] = [\bar{\varphi}_1, (\dots), \bar{\varphi}_4]$$



Edge basis functions

Volume quantities: $p=6$



Reduce
(Discretize)
 \mathcal{R}

$$[\tilde{\varphi}_i] = [\bar{\varphi}_1, (\dots), \bar{\varphi}_7]$$

Reconstruct
(Histopolate)
 \mathcal{I}

Higher dimensions: (2D case)

Higher dimensions: (2D case)

Nodal quantities:

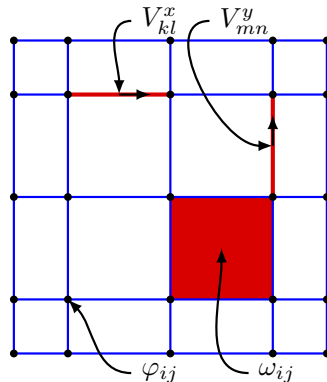
$$\varphi(x, y) \approx \tilde{\varphi}^P(x, y) = \sum_{ij} \varphi_{ij} h_i^P(x) h_j^P(y)$$

1-form quantities:

$$\begin{aligned} \vec{V}(x, y) \approx \tilde{V}^P(x, y) &= \sum_{ij} \bar{V}_{ij}^x e_i^{P-1}(x) h_j^P(y) dx \\ &+ \sum_{ij} \bar{V}_{ij}^y h_i^P(x) e_j^{P-1}(y) dy \end{aligned}$$

Volume quantities:

$$\omega(x, y) \approx \tilde{\omega}^P(x, y) = \sum_{ij} \bar{\omega}_{ij} e_i^{P-1}(x) e_j^{P-1}(y) dx dy$$



Finite volumes and Galerkin

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$$d\phi^0 = v^1$$

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$$\langle d\phi^0, \epsilon_j^1 \rangle = \langle v^1, \epsilon_j^1 \rangle$$

Finite volumes and Galerkin

$$\begin{aligned}d\phi^0 &= v^1 \\ \langle d\phi^0, \epsilon_j^1 \rangle &= \langle v^1, \epsilon_j^1 \rangle \\ \langle d\phi^0 - v^1, \epsilon_j^1 \rangle &= 0\end{aligned}$$

Finite volumes and Galerkin

$$\begin{aligned}d\phi^0 &= v^1 \\ \langle d\phi^0, \epsilon_j^1 \rangle &= \langle v^1, \epsilon_j^1 \rangle \\ \langle d\phi^0 - v^1, \epsilon_j^1 \rangle &= 0\end{aligned}$$

$$\phi^0 = \sum_i \phi_i \epsilon_i^0, \quad v^1 = \sum_i v_i \epsilon_i^1$$

Finite volumes and Galerkin

$$\langle d\phi^0 - v^1, \epsilon_j^1 \rangle = 0$$

Finite volumes and Galerkin

$$\langle d\phi^0 - v^1, \epsilon_j^1 \rangle = 0$$
$$\langle d \sum_i \phi_i \epsilon_i^0 - \sum_i v_i \epsilon_i^1, \epsilon_j^1 \rangle = 0$$

Finite volumes and Galerkin

$$\begin{aligned}\langle d\phi^0 - v^1, \epsilon_j^1 \rangle &= 0 \\ \langle d \sum_i \phi_i \epsilon_i^0 - \sum_i v_i \epsilon_i^1, \epsilon_j^1 \rangle &= 0 \\ \langle \sum_{ik} [\delta]_{ik} \phi_k \epsilon_i^1 - \sum_i v_i \epsilon_i^1, \epsilon_j^1 \rangle &= 0\end{aligned}$$

Finite volumes and Galerkin

$$\begin{aligned} \langle d\phi^0 - v^1, \epsilon_j^1 \rangle &= 0 \\ \langle d \sum_i \phi_i \epsilon_i^0 - \sum_i v_i \epsilon_i^1, \epsilon_j^1 \rangle &= 0 \\ \langle \sum_{ik} [\delta]_{ik} \phi_k \epsilon_i^1 - \sum_i v_i \epsilon_i^1, \epsilon_j^1 \rangle &= 0 \\ \sum_{ik} [\delta]_{ik} \phi_k \langle \epsilon_i^1, \epsilon_j^1 \rangle - \sum_i v_i \langle \epsilon_i^1, \epsilon_j^1 \rangle &= 0 \end{aligned}$$

Finite volumes and Galerkin

$$\begin{aligned} \langle d\phi^0 - v^1, \epsilon_j^1 \rangle &= 0 \\ \langle d \sum_i \phi_i \epsilon_i^0 - \sum_i v_i \epsilon_i^1, \epsilon_j^1 \rangle &= 0 \\ \langle \sum_{ik} [\delta]_{ik} \phi_k \epsilon_i^1 - \sum_i v_i \epsilon_i^1, \epsilon_j^1 \rangle &= 0 \\ \sum_{ik} [\delta]_{ik} \phi_k \langle \epsilon_i^1, \epsilon_j^1 \rangle - \sum_i v_i \langle \epsilon_i^1, \epsilon_j^1 \rangle &= 0 \\ \left(\sum_k [\delta]_{ik} \phi_k - v_i \right) \langle \epsilon_i^1, \epsilon_j^1 \rangle &= 0 \end{aligned}$$

Finite volumes and Galerkin

$$\begin{aligned}
 \langle d\phi^0 - v^1, \epsilon_j^1 \rangle &= 0 \\
 \langle d \sum_i \phi_i \epsilon_i^0 - \sum_i v_i \epsilon_i^1, \epsilon_j^1 \rangle &= 0 \\
 \langle \sum_{ik} [\delta]_{ik} \phi_k \epsilon_i^1 - \sum_i v_i \epsilon_i^1, \epsilon_j^1 \rangle &= 0 \\
 \sum_{ik} [\delta]_{ik} \phi_k \langle \epsilon_i^1, \epsilon_j^1 \rangle - \sum_i v_i \langle \epsilon_i^1, \epsilon_j^1 \rangle &= 0 \\
 \left(\sum_k [\delta]_{ik} \phi_k - v_i \right) \langle \epsilon_i^1, \epsilon_j^1 \rangle &= 0 \\
 \Rightarrow \sum_k [\delta]_{ik} \phi_k - v_i &= 0, \quad \forall i
 \end{aligned}$$

Conclusions

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- ▶ Algebraic topology as a natural discrete analogue of differential geometry.
- ▶ Geometrical reduction operator, \mathcal{R} , produces exact differential operators
- ▶ Reduction operator, \mathcal{R} , Whitney map, \mathcal{W} , and exterior derivative, d , have nice commuting properties.
- ▶ Edge basis functions for finite elements yield finite volume like discretizations .

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- ▶ Edge basis functions for finite elements yield finite volume like discretizations .

J. Kreeft presentation on *Higher-order discretization of Laplace operator* for application of this ideas.

Thank you